

Ph 106b Solution Set #5

- ① a) In equilibrium, \vec{E} must be 0 inside a conductor, since otherwise it would have made the charges move. Therefore, $\int \rho dV = \epsilon_0 \oint \vec{E} \cdot d\vec{s} = 0$ for an arbitrary closed surface inside the conductor. Hence, $\rho = 0$ in the interior.
- b) Since $\vec{E} = 0$, the conductor is equipotential, and so is its inner surface. Since the field in the cavity is determined from the Laplace eqn. $\Delta\phi = 0$, this means that the whole cavity is equipotential, i.e. $\vec{E} = 0$ in the cavity. The same argument will not apply outside, since the solution to $\Delta\phi = 0$ outside depends on the boundary condition at infinity.
- c) $\vec{E}_{||}$ has to be 0, since it would have made the charges flow along the surface.



\vec{E}_\perp can be determined from the Gauss law applied to a small cylinder enclosing a portion of the surface:

$$\vec{E}_\perp \cdot d\vec{s} = \frac{\rho}{\epsilon_0} \cdot d\vec{s}$$

$$E_\perp = \frac{\rho}{\epsilon_0}$$

(2) a)

From the Gauss law

$$E \cdot S = \frac{\epsilon_0}{\epsilon_0} S$$

$$E = \frac{\epsilon_0}{\epsilon_0} \quad \varphi = - \frac{\epsilon_0}{\epsilon_0} z$$

$$\Delta \varphi = \frac{\epsilon_0}{\epsilon_0} d \quad C = \frac{q}{\Delta \varphi} = \boxed{\frac{\epsilon_0 A}{d}}$$

b)

$$4\pi r^2 E = q/\epsilon_0$$

$$E = \frac{q}{4\pi r^2 \epsilon_0} \quad \varphi = \frac{q}{4\pi \epsilon_0 r} \quad \Delta \varphi = \frac{q}{4\pi \epsilon_0} \left(\frac{1}{a} - \frac{1}{b} \right)$$

$$C = \boxed{\frac{4\pi \epsilon_0 a b}{b-a}}$$

c)

$$2\pi r E = q/\epsilon_0$$

$$E = \frac{q}{2\pi r \epsilon_0} \quad \Delta \varphi = \frac{q}{2\pi \epsilon_0} \ln b/a$$

$$C = \frac{2\pi \epsilon_0}{\ln b/a}$$

d)

$$c/L = 3 \cdot 10^{-12} \text{ F/m} \rightarrow b = 6.39 \text{ mm}$$

$$c/L = 3 \cdot 10^{-12} \text{ F/m} \rightarrow b = 113 \text{ km}$$

(3)

$$\Delta \varphi = \rho/\epsilon_0 \quad \Delta \varphi' = \rho'/\epsilon_0$$

$$\oint (\rho \varphi' - \rho' \varphi) dV = \frac{1}{\epsilon_0} \int (\varphi' \Delta \varphi - \varphi \Delta \varphi') dV = \frac{1}{\epsilon_0} \int \left(\varphi' \frac{\partial \varphi}{\partial n} - \varphi \frac{\partial \varphi'}{\partial n} \right) dS =$$

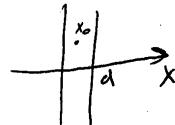
$$= - \int (\varphi' b - \varphi b') dS$$

$$\int \rho \varphi' dV + \int \varphi b' dS = \int \rho' \varphi dV + \int \varphi b' dS$$

(4) Let's choose a capacitor for our comparison problem. Then $\rho' = 0$, $\varphi' = \frac{\epsilon_0}{\epsilon_0} x$. ~~.....~~

On the other hand, $\rho = q \delta(\vec{r} - \vec{r}_0)$, where \vec{r}_0 is the position of the charge, $\varphi = 0$ on the plates, since they are grounded.

From the reciprocity theorem,



$$q \frac{\epsilon_0}{\epsilon_0} x_0 + \frac{\epsilon_0}{\epsilon_0} d \cdot Q = 0$$

where d is the distance between the plates and Q is the charge induced on the plate located at $x = d$.

$$Q = - \frac{q x_0}{d}$$